

# New hadronic states observed at BES II

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**Abstract.** The report includes the new observation of  $X(1835)$  in  $J/\psi \rightarrow \gamma\eta'\pi^+\pi^-$ , the  $\omega\phi$  threshold enhancement in  $J/\psi \rightarrow \gamma\omega\phi$ , the  $\omega\omega$  structure in  $J/\psi \rightarrow \gamma\omega\omega$  and the broad  $1^{--}$  structure at the low  $K^+K^-$  invariant-mass spectrum in  $J/\psi \rightarrow K^+K^-\pi^0$ .

**PACS.** 12.39.Mk Glueball and nonstandard multi-quark/gluon states – 13.25.Gv Decays of  $J/\psi$ ,  $\Upsilon$ , and other quarkonia

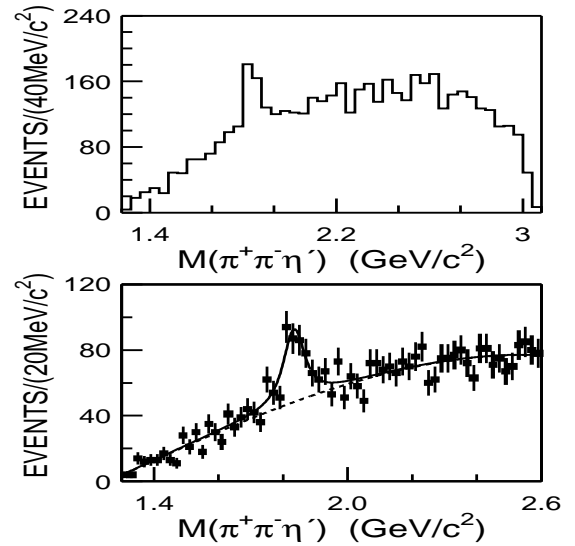
## 1 Introduction

BES [1] is a large general purpose solenoidal detector at Beijing Electron Positron Collider (BEPC). The  $5.8 \times 10^7$   $J/\psi$  events were accumulated with BES II, which provides a good laboratory for the search of new hadronic states. All the results presented here are based on this data sample.

## 2 Observation of $X(1835)$ in $J/\psi \rightarrow \gamma\eta'\pi^+\pi^-$

An anomalous enhancement near the mass threshold in the  $p\bar{p}$  invariant-mass spectrum from  $J/\psi \rightarrow \gamma p\bar{p}$  decays was reported by the BES II experiment [2]. This enhancement was fitted with a sub-threshold  $S$ -wave Breit-Wigner (BW) resonance function with a mass  $M = 1859_{-10}^{+3+5}_{-25}$  MeV/ $c^2$ , a width  $\Gamma < 30$  MeV/ $c^2$  (at the 90% C.L.) and a product branching fraction  $B(J/\psi \rightarrow \gamma X) \cdot B(X \rightarrow p\bar{p}) = (7.0 \pm 0.4_{-0.8}^{+1.9}) \times 10^{-5}$ . Among various theoretical interpretations [3–8] of the  $p\bar{p}$  mass threshold enhancement, the most intriguing one is that of a  $p\bar{p}$  bound state, sometimes called baryonium [3, 6, 9], which has been the subject of many experimental searches [10]. The baryonium interpretation of the  $p\bar{p}$  mass enhancement requires a new resonance with a mass around 1.85 GeV/ $c^2$  and it would be supported by the observation of the resonance in other decay channels.  $\pi^+\pi^-\eta'$  [5, 6] is a possible decay mode for a  $p\bar{p}$  bound state.

The analysis of  $J/\psi \rightarrow \gamma\eta'\pi^+\pi^-$  is described in detail in ref. [11]. Two decay modes of  $\eta'$  ( $\eta' \rightarrow \pi^+\pi^-\eta$ ,  $\eta \rightarrow \gamma\gamma$  and  $\eta' \rightarrow \gamma\rho$ ) are used in the analysis. The combined



**Fig. 1.** The  $\pi^+\pi^-\eta'$  invariant-mass distribution for selected events from both the  $J/\psi \rightarrow \gamma\eta'\pi^+\pi^-$  ( $\eta' \rightarrow \pi^+\pi^-\eta$ ,  $\eta \rightarrow \gamma\gamma$ ) and  $J/\psi \rightarrow \gamma\eta'\pi^+\pi^-$  ( $\eta' \rightarrow \gamma\rho$ ) analyses. The bottom panel shows the fit (solid curve) to the data (points with error bars); the dashed curve indicates the background function.

$\pi^+\pi^-\eta'$  spectrum for the two decay modes of  $\eta'$  is fitted with a BW function convolved with a Gaussian mass resolution function ( $\sigma = 13$  MeV/ $c^2$ ) to represent the  $X(1835)$  signal plus a smooth polynomial background function. The mass and width obtained from the fit (shown in the bottom panel of fig. 1) are  $M = 1833.7 \pm 6.1$  MeV/ $c^2$  and  $\Gamma = 67.7 \pm 20.3$  MeV/ $c^2$ . The signal yield from the fit is  $264 \pm 54$  events with a confidence level 45.5% ( $\chi^2/\text{d.o.f.} = 57.6/57$ ) and  $-2 \ln \mathcal{L} = 58.4$ . A fit to the mass spectrum without BW signal function returns  $-2 \ln \mathcal{L} = 126.5$ . The change

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in  $-2\ln\mathcal{L}$  with  $\Delta(\text{d.o.f}) = 3$  corresponds to a statistical significance of  $7.7\sigma$  for the signal.

Using MC-determined selection efficiencies of 3.72% and 4.85% for the  $\eta' \rightarrow \pi^+\pi^-\eta$  and  $\eta' \rightarrow \gamma\rho$  modes, respectively, we determine a product BF of

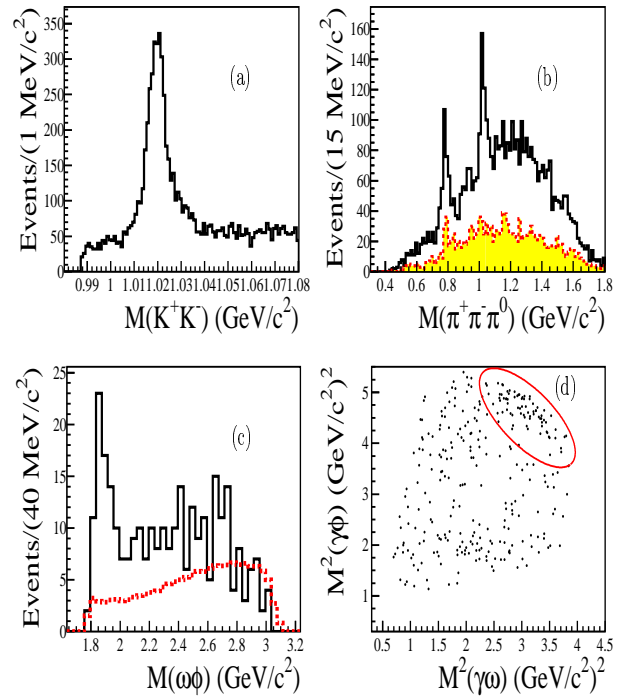
$$B(J/\psi \rightarrow \gamma X) \cdot B(X \rightarrow \pi^+\pi^-\eta') = (2.2 \pm 0.4) \times 10^{-4}.$$

The mass and width of the  $X(1835)$  are not compatible with any known meson resonance [12]. We examined the possibility that the  $X(1835)$  is responsible for the  $p\bar{p}$  mass threshold enhancement observed in radiative  $J/\psi \rightarrow \gamma p\bar{p}$  decay [2]. It has been pointed out that the  $S$ -wave BW function used for the fit in ref. [2] should be modified to include the effect of final-state interactions (FSI) on the shape of the  $p\bar{p}$  mass spectrum [7,8]. Redoing the  $S$ -wave BW fit to the  $p\bar{p}$  invariant-mass spectrum in ref. [2] including the zero isospin,  $S$ -wave FSI factor of ref. [8], yields a mass  $M = 1831 \pm 7 \text{ MeV}/c^2$  and a width  $\Gamma < 153 \text{ MeV}/c^2$  (at the 90% C.L.); these values are in good agreement with the mass and width of  $X(1835)$  reported here. Moreover, according to ref. [6], the  $\pi\pi\eta'$  mode is expected to be strong for a  $p\bar{p}$  bound state. Thus, the  $X(1835)$ -resonance is a prime candidate for the source of the  $p\bar{p}$  mass threshold enhancement in  $J/\psi \rightarrow \gamma p\bar{p}$  process. In this case, the  $J^{PC}$  and  $I^G$  of the  $X(1835)$  could only be  $0^{++}$  and  $0^+$ , which can be tested in future experiments. Also in this context, the relative  $p\bar{p}$  decay strength is quite strong:  $B(X \rightarrow p\bar{p})/B(X \rightarrow \pi^+\pi^-\eta') \sim 1/3$  (The product BF determined from the fit that includes FSI effects on the  $p\bar{p}$  mass spectrum is within the systematic errors of the result report in ref. [2]). Since decays to  $p\bar{p}$  are kinematically allowed only for a small portion of the high-mass tail of the resonance and have very limited phase space, the large  $p\bar{p}$  branching fraction implies an unusually strong coupling to  $p\bar{p}$ , as expected for a  $p\bar{p}$  bound state [9,13]. However, other possible interpretations of the  $X(1835)$  that have no relation to the  $p\bar{p}$  mass threshold enhancement are not excluded.

### 3 Observation of the $\omega\phi$ threshold enhancement in $J/\psi \rightarrow \gamma\omega\phi$

Systems of two vector particles have been intensively studied for signatures of gluonic bound states. The radiative  $J/\psi$  decay  $J/\psi \rightarrow \gamma\omega\phi$  is a double Okubo-Zweig-Iizuka (OZI) suppressed process, the measurement of this decay and the search for possible resonance states will provide useful information on two vector-meson systems. MARK III Collaboration [14] studied  $J/\psi \rightarrow \gamma\omega\phi$  decays, but did not find clear structures in the  $\omega\phi$  invariant-mass spectrum. The final states of  $\omega\phi$  are also observed in photon-photon collisions by ARGUS [15,16] experiment and the cross-sections were measured [16].

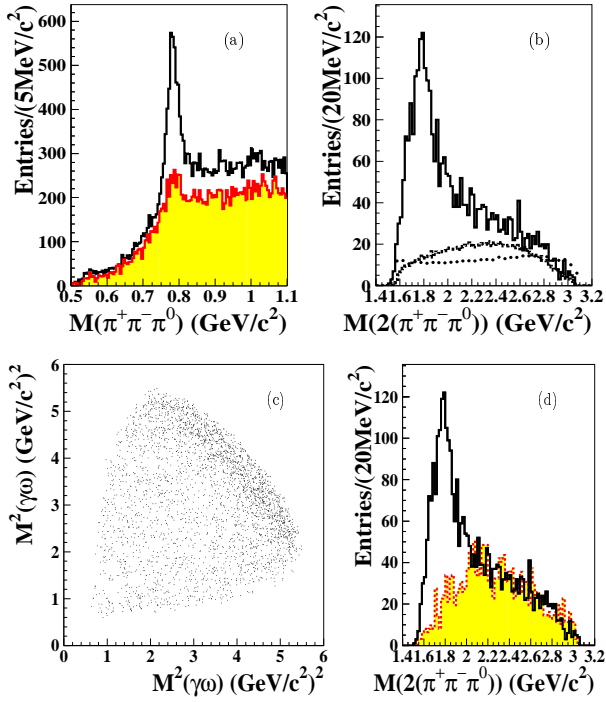
The detailed analysis of  $J/\psi \rightarrow \gamma\omega\phi$  is described in ref. [17]. The decay modes of  $\omega \rightarrow \pi^+\pi^-\pi^0$  and  $\phi \rightarrow K^+K^-$  are used in the analysis. Since the decays of  $J/\psi \rightarrow \omega\phi$  and  $J/\psi \rightarrow \omega\phi\pi^0$  are forbidden by  $C$  invariance, the 294 observed  $\omega\phi$  events present direct evidence



**Fig. 2.** (a) The  $K^+K^-$  invariant-mass distribution. (b) The  $\pi^+\pi^-\pi^0$  invariant-mass distribution; the open histogram is for candidate events with  $m_{K^+K^-}$  being in the  $\phi$  range, and the shaded histogram is for events with  $m_{K^+K^-}$  being in the  $\phi$  sideband region. (c) The  $K^+K^-\pi^+\pi^-\pi^0$  invariant-mass distribution for the  $J/\psi \rightarrow \gamma\omega\phi$  candidate events. The dashed curve indicates the acceptance varying with the  $\omega\phi$  invariant mass. (d) Dalitz plot.

for the radiative  $J/\psi \rightarrow \gamma\omega\phi$  decay. The histogram in fig. 2(c) shows the  $K^+K^-\pi^+\pi^-\pi^0$  invariant-mass distribution for events with  $|m_{K^+K^-} - m_\phi| < 15 \text{ MeV}/c^2$  and  $|m_{\pi^+\pi^-\pi^0} - m_\omega| < 30 \text{ MeV}/c^2$ , and a structure peaked near the  $\omega\phi$  threshold is observed. The dashed curve in the figure indicates how the acceptance varies with invariant mass. The peak is also evident as a diagonal band along the upper right-hand edge of the Dalitz plot in fig. 2(d). No evidence of an enhancement near the  $\omega\phi$  threshold is observed from the sideband events of  $\omega$  and  $\phi$ . And the study from inclusive and exclusive MC shows that the  $\omega\phi$  threshold enhancement is not from backgrounds.

The significance of the  $\omega\phi$  threshold enhancement is more than  $10\sigma$ . From a partial-wave analysis with covariant helicity coupling amplitudes, a spin-parity of  $X = 0^{++}$  with an  $S$ -wave  $\omega\phi$  system is favored. The mass and width of the enhancement are determined to be  $M = 1812^{+19}_{-26} \text{ (stat)} \pm 18 \text{ (syst)} \text{ MeV}/c^2$  and  $\Gamma = 105 \pm 20 \text{ (stat)} \pm 28 \text{ (syst)} \text{ MeV}/c^2$ , and the product branching fraction is  $B(J/\psi \rightarrow \gamma X) \cdot B(X \rightarrow \omega\phi) = (2.61 \pm 0.27 \text{ (stat)} \pm 0.65 \text{ (syst)}) \times 10^{-4}$ . The mass and width of this state are not compatible with any known scalars listed in the Particle Data Group (PDG) [12]. It could be an unconventional state [18–22]. However, more statistics and further studies are needed to clarify this.

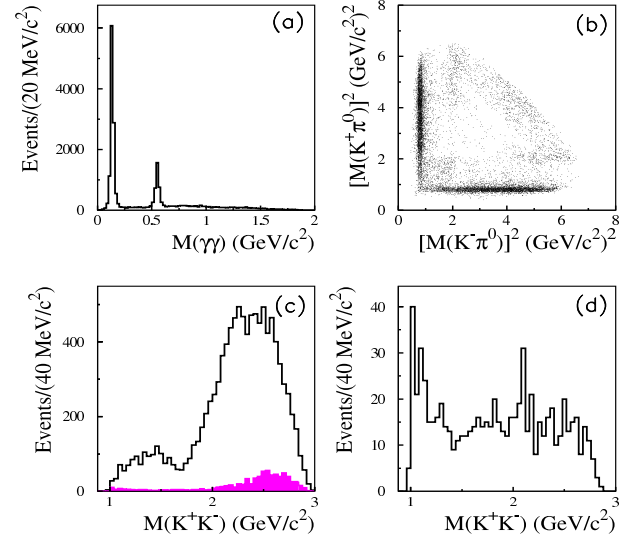


**Fig. 3.** (a) The  $\pi^+\pi^-\pi^0$  mass distribution for the best  $\omega\omega$  combination after the requirement that the invariant mass of the other  $\pi^+\pi^-\pi^0$  is inside the  $\omega$  region (open histogram), defined by  $|m_{\pi^+\pi^-\pi^0} - m_\omega| < 40 \text{ MeV}/c^2$ , or in the sideband range (shaded histogram), defined by  $40 \text{ MeV}/c^2 < |m_{\pi^+\pi^-\pi^0} - m_\omega| < 80 \text{ MeV}/c^2$ . (b) The  $2(\pi^+\pi^-\pi^0)$  invariant-mass distribution for candidate events. The dashed curve is the phase space invariant-mass distribution, and the dotted curve shows the acceptance versus the  $\omega\omega$  invariant mass. (c) The Dalitz plot. (d) The  $2(\pi^+\pi^-\pi^0)$  invariant mass of the inclusive Monte Carlo sample (shaded histogram).

#### 4 Pseudoscalar production at the $\omega\omega$ threshold in $J/\psi \rightarrow \gamma\omega\omega$

The  $\eta(1760)$  was reported by the MARK III Collaboration in  $J/\psi$  radiative decays and was found to decay to  $\omega\omega$  [23] and  $\rho\rho$  [24]. It was also observed by the DM2 [25, 26] and BES I [27]. The  $\eta(1760)$  was suggested to be a  $3^1S_0$  pseudoscalar  $q\bar{q}$  meson, but some authors suggest a mixture of glueball and  $q\bar{q}$  or a hybrid [28, 29]. In our analysis of  $J/\psi \rightarrow \gamma\omega\omega$  ( $\omega \rightarrow \pi^+\pi^-\pi^0$ ) [30], the presence of a signal around  $1.76 \text{ GeV}/c^2$  and its pseudoscalar character are confirmed, and the mass, width, and branching fraction are measured by partial-wave analysis.

The histogram of fig. 3(b) shows the  $2(\pi^+\pi^-\pi^0)$  invariant-mass distribution of events with both  $\pi^+\pi^-\pi^0$  masses within the  $\omega$  range ( $|m_{\pi^+\pi^-\pi^0} - m_\omega| < 40 \text{ MeV}/c^2$ ). There are 3046 events with a clear peak at  $1.76 \text{ GeV}/c^2$ . The phase space invariant-mass distribution and the acceptance versus  $\omega\omega$  invariant mass are also shown in the figure. The corresponding Dalitz plot is shown in fig. 3(c). The shaded histogram of fig. 3(d) shows the  $2(\pi^+\pi^-\pi^0)$  invariant-mass distribution of the inclusive sample. There is no peak at the  $\omega\omega$  mass threshold



**Fig. 4.** (a) The  $\gamma\gamma$  invariant-mass distribution. (b) The Dalitz plot for  $K^+K^-\pi^0$  candidate events. (c) The  $K^+K^-$  invariant-mass distribution for  $K^+K^-\pi^0$  candidate events; the solid histogram is data and the shaded histogram is the background (normalized to data). (d) The  $K^+K^-$  invariant-mass distribution for the  $\pi^0$  mass sideband events (not normalized).

in the invariant-mass distribution at around  $1.76 \text{ GeV}/c^2$ . The fitting of fig. 3(d) yields  $1441 \pm 50$  background events within the  $\omega\omega$  invariant-mass range from  $1.6 \text{ GeV}/c^2$  to  $2.8 \text{ GeV}/c^2$ .

The partial-wave analysis shows that the structure around  $1.76 \text{ GeV}/c^2$  in the  $\omega\omega$  invariant-mass spectrum is predominantly pseudoscalar, with small contributions from  $f_0(1710)$ ,  $f_2(1640)$ , and  $f_2(1910)$ . The mass of the pseudoscalar is  $M = 1744 \pm 10$  (stat)  $\pm 15$  (syst)  $\text{MeV}/c^2$ , the width  $\Gamma = 244_{-21}^{+24}$  (stat)  $\pm 25$  (syst)  $\text{MeV}/c^2$ , and the product branching fraction is  $\text{Br}(J/\psi \rightarrow \gamma\eta(1760)) \cdot \text{Br}(\eta(1760) \rightarrow \omega\omega) = (1.98 \pm 0.08$  (stat)  $\pm 0.32$  (syst))  $\times 10^{-3}$ .

#### 5 Observation of a broad $1^{--}$ resonant structure in the $K^+K^-$ mass spectrum in $J/\psi \rightarrow K^+K^-\pi^0$

A broad peak is observed at low  $K^+K^-$  invariant mass in  $J/\psi \rightarrow K^+K^-\pi^0$  decays, detailed analysis is described in ref. [31]. The Dalitz plot for the selected 10631 events is shown in fig. 4(b), where a broad  $K^+K^-$  band is evident in addition to the  $K^*(892)$  and  $K^*(1410)$  signals. This band corresponds to the broad peak observed around  $1.5 \text{ GeV}/c^2$  in the  $K^+K^-$  invariant mass projection shown in fig. 4(c).

A partial-wave analysis shows that the  $J^{PC}$  of this structure is  $1^{--}$ . Its pole position is determined to be  $(1576_{-55-91}^{+49+98}) \text{ MeV}/c^2 - i(409_{-12-67}^{+11+32}) \text{ MeV}/c^2$ , and the branching ratio is  $B(J/\psi \rightarrow X\pi^0) \cdot B(X \rightarrow K^+K^-) = (8.5 \pm 0.6_{-3.6}^{+2.7}) \times 10^{-4}$ , where the first errors are statistical

and the second are systematic. These parameters are not compatible with any known meson resonances [12].

To understand the nature of the broad  $1^{--}$  peak, it is important to search for a similar structure in  $J/\psi \rightarrow K_S^0 K^\pm \pi^\mp$  decays to determine its isospin. It is also intriguing to search for  $K^* K, K \bar{K} \pi$  decay modes. In the mass region of the  $X$ , there are several other  $1^{--}$  states, such as the  $\rho(1450)$  and  $\rho(1700)$ , but the width of the  $X$  is much broader than the widths of these other mesons. This may be an indication that the  $X$  has a different nature than these other mesons. For example, very broad widths are expected for multiquark states [32].

## 6 Summary

Based on  $5.8 \times 10^7$   $J/\psi$  events accumulated at the BES II detector, the  $X(1835)$  in  $J/\psi \rightarrow \gamma \eta' \pi^+ \pi^-$ , the  $\omega \phi$  threshold enhancement in  $J/\psi \rightarrow \gamma \omega \phi$  and a broad  $1^{--}$  resonant structure around  $1.5 \text{ GeV}/c^2$  in the  $K^+ K^-$  mass spectrum in  $J/\psi \rightarrow K^+ K^- \pi^0$  are observed. And partial-wave analysis shows that  $\eta(1760)$  is dominant in the  $\omega \omega$  in the invariant-mass spectrum in  $J/\psi \rightarrow \gamma \omega \omega$ .

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